

Cluster algebras for spinor helicity and momentum twistor varieties

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In quantum field theory physicist probability of particle interactions are computed via scattering amplitudes.

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The *symbol* of an iterated integral captures information on the integration kernels while disregarding coefficients, e.g.

$$Li_2(x) = - \int_0^x d \log(y) \int_0^z d \log(1 - z)$$

has symbol $\mathcal{S}(Li_2(x)) = -(1 - x) \otimes x, 1 - x$ and x are called the *letters*.

Motivation: cluster bootstrap

In the toy model *planar* $\mathcal{N} = 4$ *super Yang-Mills* symbol letters of the amplitudes for $n = 6, 7$ particles are *cluster variables* Grassmannian $\text{Gr}_{4,n}$.

This has led to a *bootstrap* for the amplitude: it is the unique generalized polylogarithmic function with symbol letters given by cluster variables.

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Question: Can this be extended beyond $\mathcal{N} = 4$ sYM?

- 1 What are the relevant kinematic spaces?
- 2 Are there cluster algebras?

Overview

- 1 Partial flag varieties
- 2 Spinor helicity and momentum twistor varieties
- 3 Cluster algebras
- 4 Symbol alphabet for the spinor helicity variety
- 5 Summary and open questions

§1 Partial flag varieties

Consider $1 \leq d_1 < d_2 < n$ and define the (two step) *partial flag variety*

$$\mathcal{F}_{d_1, d_2; n} := \{0 \in V_1 \subsetneq V_2 \subsetneq \mathbb{C}^n : \dim_{\mathbb{C}} V_i = d_i\}$$

Associate to $\mathcal{V} \in \mathcal{F}_{d_1, d_2; n}$ a matrix $M_{\mathcal{V}} = (m_{ij}) \in \mathbb{C}^{d_2 \times n}$ such that V_i is generated by the first d_i rows of M .

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$\{i_1, \dots, i_{d_j}\} \subset \{1, \dots, n\}$, define the *Plücker coordinate*

$$P_{i_1, \dots, i_{d_j}}(\mathcal{V}) := \det(m_{ab})_{1 \leq a \leq d_j, b \in \{i_1, \dots, i_{d_j}\}}.$$

Example: $M_{\mathcal{V}} = \begin{bmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 1 & 0 & 0 & -3 \end{bmatrix}$, so that $P_{134}(\mathcal{V}) = \det \left(\begin{bmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 1 & 0 & -3 \end{bmatrix} \right) = -6$.

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The $P_{i_1, \dots, i_{d_j}}$ satisfy *Plücker relations* (quadratic polynomials) and determine an embedding

$$\mathcal{F}_{d_1, d_2; n} \hookrightarrow \mathbb{P}^{\binom{n}{d_1}-1} \times \mathbb{P}^{\binom{n}{d_2}-1}$$

Let $\mathbb{C}[\mathcal{F}_{d_1, d_2; n}]$ denote the homogeneous coordinate ring.

§2 Example 1: Particles and helicity spinors

Minkowski space \mathbb{R}^4 with inner product

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with $\lambda = \frac{1}{\sqrt{p_0 + p_3}} \begin{pmatrix} p_0 + p_3 \\ p_1 + ip_2 \end{pmatrix}$, $\tilde{\lambda} = \frac{1}{\sqrt{p_0 + p_3}} \begin{pmatrix} p_0 + p_3 \\ p_1 - ip_2 \end{pmatrix} \in \mathbb{C}^2$ we have

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The mass m of p satisfies $p^2 := p \cdot p = m^2$. If p is *lightlike* $m = 0$, we have

$$\det(\sigma(p)) = p_0^2 - (p_1^2 + p_2^2 + p_3^2) = p^2 = 0.$$

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satisfying *Schouten identities* and *momentum conservation*

$$0 = P_{ij} P_{kl} - P_{ik} P_{jl} + P_{il} P_{jk} = \tilde{P}_{ij} \tilde{P}_{kl} - \tilde{P}_{ik} \tilde{P}_{jl} + \tilde{P}_{il} \tilde{P}_{jk}$$

$$0 = \sum_{s=1}^n P_{is} \tilde{P}_{sj} \quad (\Leftrightarrow \Lambda \tilde{\Lambda}^T = 0)$$

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Consider the map

$$\langle ij \rangle \mapsto P_{ij}, \quad \text{and} \quad [ij] \mapsto (-1)^{i+j-1} P_{[n]-ij}$$

where $[n] - ij := \{1, \dots, n\} - \{i, j\}$.

¹(*) in [Y.El Mazzouz, A.Pfister, and B.Sturmfels. 2406.17331]

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$$\begin{aligned} P_{ij}P_{kl} - P_{ik}P_{jl} + P_{il}P_{jk} &= 0 \\ P_{[n]-ij}P_{[n]-kl} - P_{[n]-ik}P_{[n]-jl} + P_{[n]-il}P_{[n]-jk} &= 0 \\ \sum_{s=1}^n (-1)^{s+j-1} P_{is}P_{[n]-js} &= 0 \end{aligned}$$

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As $\dim \mathcal{SH}_n \stackrel{(*)}{=} 4(n-3) = \dim \mathcal{F}_{2,n-2;n}$ the map induces an isomorphism (identification) between the spinor helicity variety \mathcal{SH}_n and the **partial flag variety** $\mathcal{F}_{2,n-2;n}$.¹

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§2 Example 2: *Momentum twistor variety*

For dual conformal invariant scattering consider momentum twistors $Z_1, \dots, Z_n \in \mathbb{C}P^3$. The system is parametrized by

$$\langle ijkl \rangle := \det(Z_i Z_j Z_k Z_l),$$

for $1 \leq i < j < k < l \leq n$ satisfying determinantal identities (Plücker relations). Identifying $\langle ijkl \rangle \mapsto P_{ijkl}$ the same equations determine the **Grassmannian** $Gr_{4,n}$ (up to scaling, i.e. $(\mathbb{C}^*)^n$ -action).

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In the non-dual conformal invariant setting introduce $Z_{n+1}, Z_{n+2} \in \mathbb{C}P^3$ to break symmetry and parametrize by

$$\langle ijkl \rangle := \det(Z_i Z_j Z_k Z_l), \quad \langle ij \rangle := \det(Z_i Z_j Z_{n+1} Z_{n+2}),$$

for $1 \leq i < j < k < l \leq n$ resp. $1 \leq i < j \leq n$.

§2 Example 2: *Momentum twistor variety*

This motivates us to define the *momentum twistor variety* as the subvariety of $\mathbb{P}^{\binom{n}{2}-1} \times \mathbb{P}^{\binom{n}{4}-1}$ determined by the vanishing of determinantal identities among the $\langle ijkl \rangle$ and $\langle ij \rangle$.

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Consider the map

$$\langle ijkl \rangle \mapsto P_{ijkl} \quad \text{and} \quad \langle ij \rangle \mapsto P_{ij}$$

So the momentum twistor variety \mathcal{MT}_n is the *partial flag variety* $\mathcal{F}_{2,4;n}$.

§3 Cluster algebras [S. Fomin, A. Zelevinsky math/0104151]

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Example: The cluster algebras for $1 \rightarrow 2$ is recursively generated by $\{x_1, x_2\}$ and the mutation rule $x_{i+1}x_{i-1} = x_i + 1$ for generating clusters $\{x_i, x_{i+1}\}$.

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$$x_1, x_2, x_3 = \frac{x_2 + 1}{x_1}, x_4 = \frac{x_1 + x_2 + 1}{x_1 x_2}, x_5 = \frac{x_1 + 1}{x_2}, x_6 = x_1, x_7 = x_2, \dots$$

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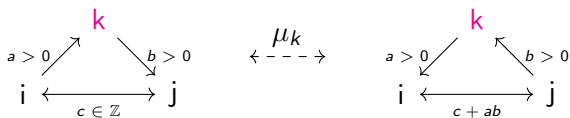
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The cluster algebras is

$$A_{1 \rightarrow 2} = \left\langle x_1, x_2, \frac{x_2 + 1}{x_1}, \frac{x_1 + x_2 + 1}{x_1 x_2}, \frac{x_1 + 1}{x_2} \right\rangle \subset \mathbb{C}(x_1, x_2)$$

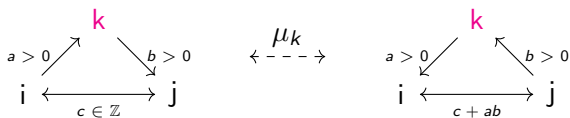
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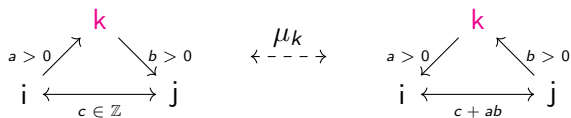


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where

$$x'_k = \frac{1}{x_k} \left(\prod_{i \rightarrow k} x_i + \prod_{k \rightarrow j} x_j \right)$$

A pair $s = (Q, \underline{x} = (x_1, \dots, x_n))$ is called a *seed*.

III. Cluster algebra

Take away: to define a cluster algebra, I need

- 1 a quiver (directed graph, no loops, no 2-cycles), and
- 2 a tuple of algebraically independent elements (one for each vertex) in a field of rational functions, and
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Recall the Plücker embedding

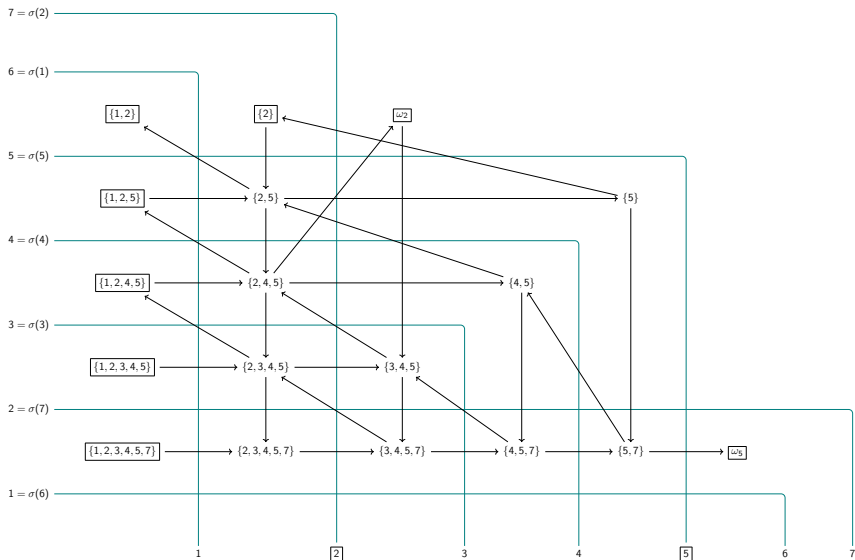
$$\mathcal{F}_{d_1, d_2; n} \hookrightarrow \mathbb{P}^{\binom{n}{d_1}-1} \times \mathbb{P}^{\binom{n}{d_2}-1}$$

with homogeneous coordinate ring $\mathbb{C}[\mathcal{F}_{d_1, d_2; n}]$

Theorem (C. Geiss, B. Leclerc, J. Schröer, Ann. Inst. Fourier 2008)

The homogeneous coordinate ring $\mathbb{C}[\mathcal{F}_{d_1, d_2; n}]$ is a cluster algebra.

III. Example $\mathcal{F}_{2,5;7}$



III. Example $\mathcal{F}_{2,5;7}$: cluster variables

Index sets are all of form $[i_j, d_j] \cup [i_{j+1}, d_{j+1}]$ and define right bound minors of a $n \times n$ unipotent upper triangular matrix.

Translating to Plücker coordinates yields the cluster variables:

$$\Delta_{[i_j, d_j] \cup [i_{j+1}, d_{j+1}]} \stackrel{\text{Laplace}}{=} \sum_{J \in \binom{[\ell, n]}{d_j - i_{j+1}}, J' = [\ell, n] \setminus J} (-1)^{\Sigma(i_j, d_j, J)} P_{[i_j - 1] \cup J} P_{[i_{j+1} - 1] \cup J'}$$

where $\ell = n - d_j - d_{j+1} + i_j + i_{j+1} - 1$.

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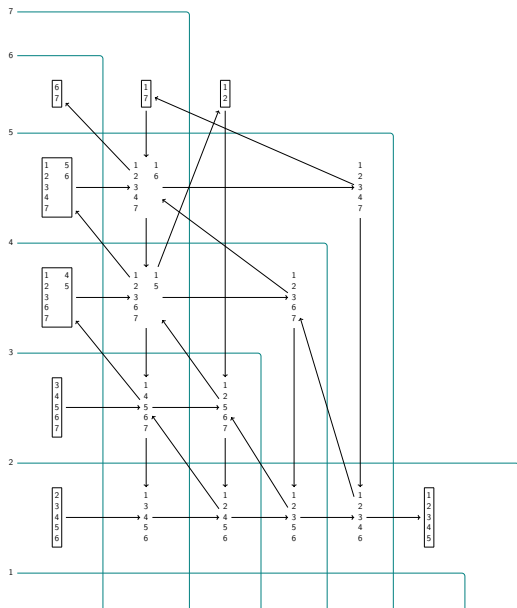
Translating to Plücker coordinates yields the cluster variables:

$$\Delta_{[i_j, d_j] \cup [i_{j+1}, d_{j+1}]} \stackrel{\text{Laplace}}{=} \sum_{J \in \binom{[\ell, n]}{d_j - i_{j+1}}, J' = [\ell, n] \setminus J} (-1)^{\Sigma(i_j, d_j, J)} P_{[i_j - 1] \cup J} P_{[i_{j+1} - 1] \cup J'}$$

where $\ell = n - d_j - d_{j+1} + i_j + i_{j+1} - 1$. Associate a two-column tableaux filled index sets of maximal term on RHS (tableaux of a cluster variable as an element of *Lusztig's dual canonical basis* [Lusztig '90, Chari–Pressley '91, Hernandez–Leclerc '10, Li '24])

Denote by $\mathcal{T}_{d_1, d_2; n}$ (resp. $\mathcal{T}_{d; n}$) the set of semistandard Young tableaux with columns of size d_1 or d_2 (resp. d), filled with $\{1, \dots, n\}$.

III. Example $\mathcal{F}_{2,5;7}$



§3 Relation between the cluster structures

There's a combinatorial map on tableaux:

$$\varphi : \mathcal{T}_{d_1, d_2; n} \hookrightarrow \mathcal{T}_{d_2; n+d_2-d_1}$$

obtained by "filling up" d_1 -columns to d_2 columns by adding $n+1, \dots, n+d_2-d_1$.

Example: $\varphi : \mathcal{T}_{2,4;6} \hookrightarrow \mathcal{T}_{4,8}$, for $a < b$ and $i < j < k < l$ in $\{1, \dots, 6\}$

$$\begin{array}{c} a \\ b \end{array} \mapsto \begin{array}{c} a \\ b \\ 7 \\ 8 \end{array}, \quad \begin{array}{c} i \\ j \\ k \\ l \end{array} \mapsto \begin{array}{c} i \\ j \\ k \\ l \end{array}$$

Theorem (L.B.-Jianrong Li, arxiv:2408.14956)

Let $s_0^F = (Q^F, \underline{x}^F)$ denote the initial seed for $\mathcal{F}_{2, n-2; n}$ (resp. $\mathcal{F}_{2,4; n}$). Then there exists a seed $s^{Gr} = (Q^{Gr}, \underline{x}^{Gr})$ for $Gr_{n-2; 2n-4}$ (resp. $Gr_{4; n+2}$) so that

- 1 Q^F is a full subquiver in Q^{Gr} obtained by freezing $n-5$ (resp. one) vertices and eliminating $(n-3)(n-5)$ (resp. 3) vertices
- 2 $\varphi(\underline{x}^F) \subset \underline{x}^{Gr}$.

IV. Example: $n = 6$ s_0^{Gr} and s^{Gr}

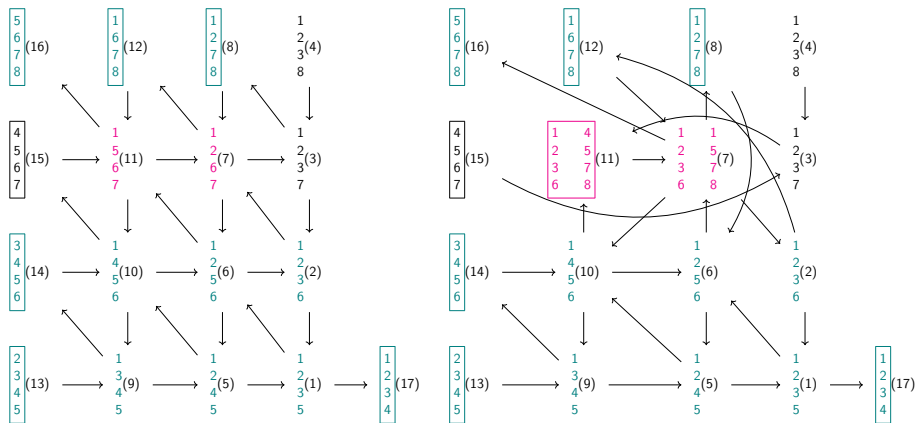


Figure: LHS: The initial seed for the Grassmannian $Gr_{4;8}$. Mutation at (7), then (11) followed by freezing (11) yields the full subquiver depicted on the RHS on all vertices but (3),(4),(15) that coincides with the initial seed for $\mathcal{F}_{2,4;6}$.

IV. Example: $n = 8$ initial seed s_0^{Gr}

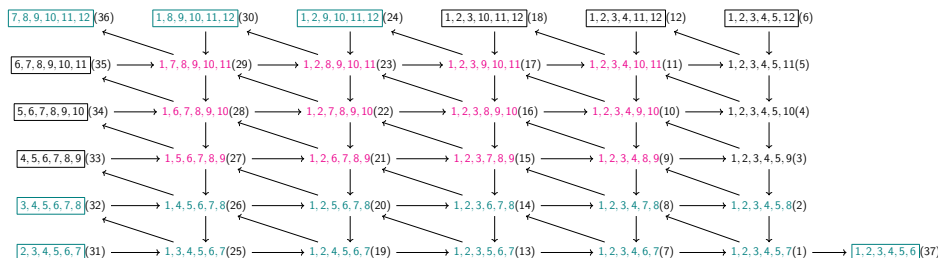


Figure: The initial seed for the Grassmannian $Gr_{6,12}$. The central 3×4 -grid from vertex (9) to (29) marked in **magenta** is the mutation grid. After the mutation sequence we freeze vertices (29), (16) and (10). Vertices marked in **teal** are those that already correspond to initial cluster variables for $\mathcal{F}_{2,6;8}$.

IV. Example: $n = 8$ mutated seed s^{Gr}

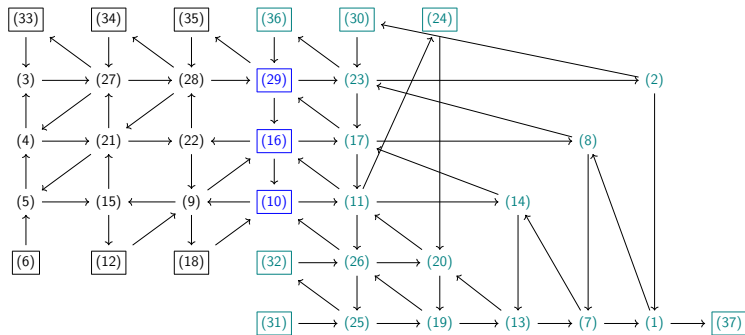
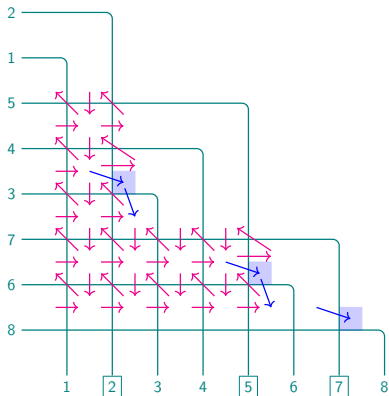


Figure: The quiver Q^{Gr} after freezing vertices 10, 16 and 29. All colored vertices on the RHS yield the subquiver corresponding to the initial quiver Q^F for $\mathcal{F}_{2,6;8}$.

What's next?

- \mathcal{SH}_n cluster structure relevant for the symbol alphabet for $n = 5$ [L.B.-Drummond-Glew], $n \geq 6$?
- Embeddings generalize for arbitrary partial flag varieties, conjecturally also compatible with cluster structure
- Application to cluster adjacency conjecture for the momentum amplituhedron
- Scattering with massive particles can be obtained from \mathcal{MT}_n by imposing vanishing certain conditions (in terms of differential operators). What is the kinematic spaces? Does it have cluster structure?

Thank you!



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